

6. Quantum Theory - Introduction

6.1 Failures of classical physics

• Black body radiation

According to the equipartition theorem all the oscillators of the field share equally the energy supplied by the heated walls \Rightarrow even the highest frequencies \Rightarrow ultraviolet catastrophe

\Rightarrow Planck:

oscillators are excited only if they can acquire an energy of at least $h\nu$
quantization of energy: $E = n h \nu$

Planck's constant
 $6.62608 \cdot 10^{-34} \text{ Js}$

• Heat capacities

Dulong and Petit's law (classical)

each atom can oscillate in three dimensions, the average energy of each atom is $3kT$, for N atoms: $3NkT$

$$U_m = 3 N_A k T = 3 R T$$

$$C_{V,m} = \left. \frac{\partial U_m}{\partial T} \right|_V = 3 R$$

fails as ~~as~~ as $T \rightarrow 0$

Einstein:

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each atom oscillated about its equilibrium position with a single frequency and the energy of oscillation is confined to discrete values $n h \nu$

$$\Rightarrow U_{\text{m}} = \frac{3 N_A h \nu}{e^{\frac{h \nu}{k T}} - 1} \Rightarrow C_{V, \text{m}} = 3 R f^2$$

$$f = \frac{\theta_E}{T} \left(\frac{e^{\theta_E / 2T}}{e^{\theta_E / T} - 1} \right)$$

Einstein temperature

$$\theta = \frac{h \nu}{k}$$

for $T \gg \theta_E \Rightarrow$ classical result

$$\text{for } T \ll \theta_E: f \approx \frac{\theta_E}{T} \left(\frac{e^{\theta_E / 2T}}{e^{\theta_E / T}} \right) = \frac{\theta_E}{T} e^{-\theta_E / 2T}$$

$$f \rightarrow 0 \text{ for } T \rightarrow 0$$

- at low temperatures only a few oscillators possess enough energy to oscillate significantly
- at higher temperatures there is enough energy available for all all oscillators to become active
 \Rightarrow classical value

Debye: range of discrete frequencies

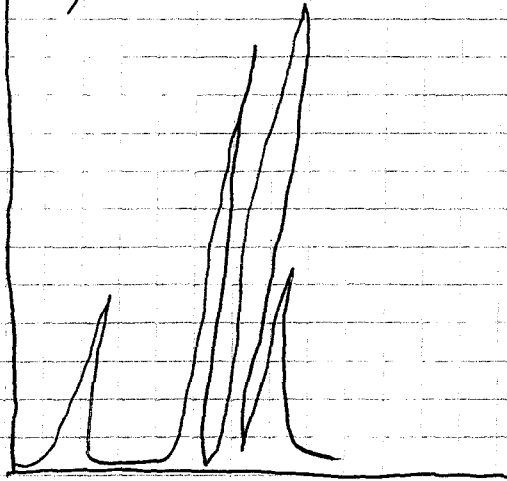
$$C_{V, \text{m}} = 3 R f \quad f = 3 \left(\frac{T}{\theta_D} \right)^3 \int_0^{\theta_D / T} \frac{x^4 e^x}{(e^x - 1)^2} dx$$
$$\theta_D = \frac{h \nu_D}{k}$$

Atomic and Molecular spectra

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atomic spectra:

emission intensity
of excited atom



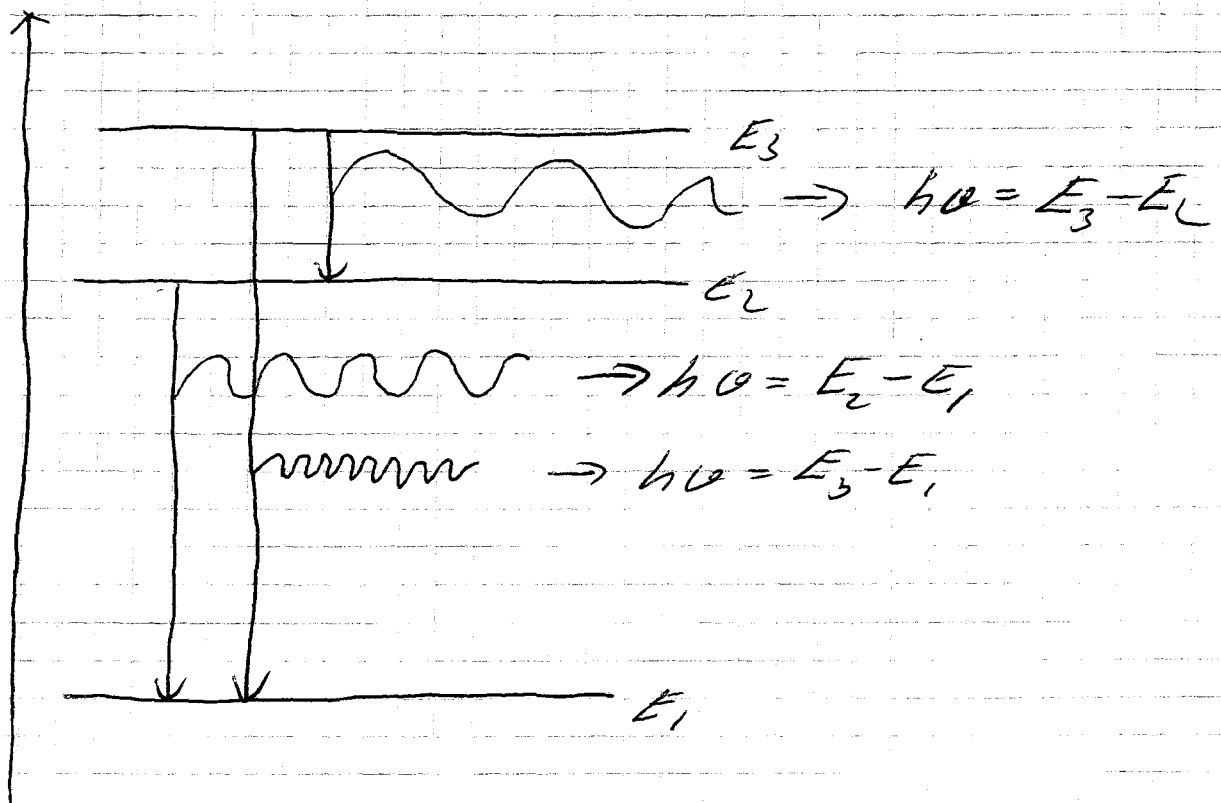
⇒ sharp lines

in absorption and emission

⇒ energy of atoms is

also confined to discrete
values

⇒ energy can be emitted or absorbed only in discrete
amounts



Bohr Radius a_0 :

potentielle Energie des Elektrons hängt allein von seinem Abstand vom Kern ab

$$U(r) = - \frac{Z \cdot e^2}{(4\pi\epsilon_0) r}$$

$$a_0 = \frac{\hbar^2 4\pi\epsilon_0}{m_e \cdot e^2}$$

$$R_n = n a_0 \quad n = 1, 2, \dots$$

$$E_n = - \frac{e^4 m_e Z^2}{(4\pi\epsilon_0)^2 2\hbar^2 n^2}$$

• Photoelectric effect

ejection of an electron when it is involved in a collision with a photon

$$\frac{1}{2} m_e v^2 = h\nu - \phi$$

↑ kinetic energy ↓ photon ← work function

• Electron diffraction

⇒ electron microscopy

de Broglie wave length:

$$\lambda = \frac{h}{p}$$

The Schrödinger equation: Ψ : wavefunction

1-dim
$$-\frac{\hbar^2}{2m} \frac{d^2 \Psi}{dx^2} + V(x) \Psi = E \Psi$$

|
potential energy

$\hbar = \frac{h}{2\pi}$

3-dim
$$-\frac{\hbar^2}{2m} \nabla^2 \Psi + V \Psi = E \Psi$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

spherical
$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \Lambda^2$$

$$\Lambda^2 = \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta}$$

$$H \Psi = E \Psi$$

Hamiltonian operator =
$$-\frac{\hbar^2}{2m} \nabla^2 + V$$

time-dependent:

$$H \Psi = i \hbar \frac{\partial \Psi}{\partial t}$$

e.g. $V(x) = V$

$$\frac{d^2 \Psi}{dx^2} = -\frac{2m}{\hbar^2} (E - V) \Psi$$

$$\Rightarrow \Psi = e^{ikx} \quad \text{with } k = \left(\frac{2m(E-V)}{\hbar^2} \right)^{1/2}$$

$$E - V = E_{\text{kin}} \Rightarrow E_{\text{kin}} = \frac{k^2 \hbar^2}{2m} = \frac{p^2}{2m}$$

$$\Rightarrow p = \hbar k = \frac{h}{\lambda} \quad \text{with } \lambda = \frac{2\pi}{k}$$

Born interpretation: ψ probability amplitude $|\psi|^2$ probability density

If the wavefunction of a particle has the value ψ at some point r then the probability of finding the particle in an infinitesimal volume $d\tau = dx dy dz$ at that point is proportional to $|\psi|^2 d\tau$

normalization $\int |\psi|^2 d\tau = 1$

ψ : continuous, continuous slope, single-valued, finite, cannot be zero everywhere

Eigenvalues and eigenfunctions