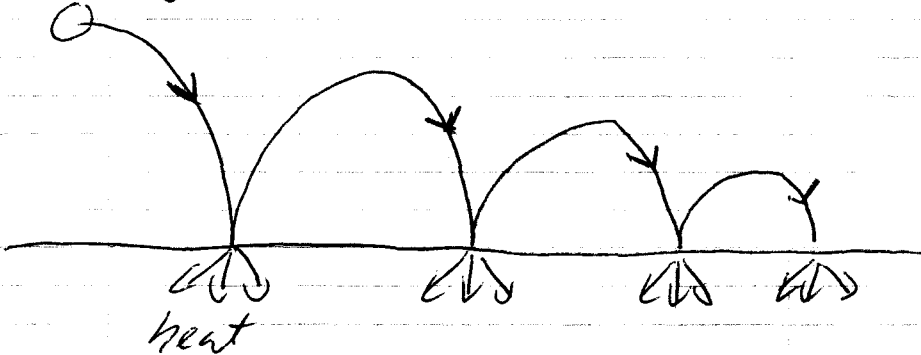


## 1.5. Entropy

bouncing ball:



no heat flow back into the ball!

⇒ spontaneous change:

the direction of change that leads to more disorderly disposal of the total energy of the isolated system!

⇒ less information

⇒ more possible states or configurations

⇒ self-assembly is not possible in an isolated system!

### Entropy $S$

The entropy of an isolated system in the course of a spontaneous change:

$$\Delta S_{\text{tot}} > 0$$

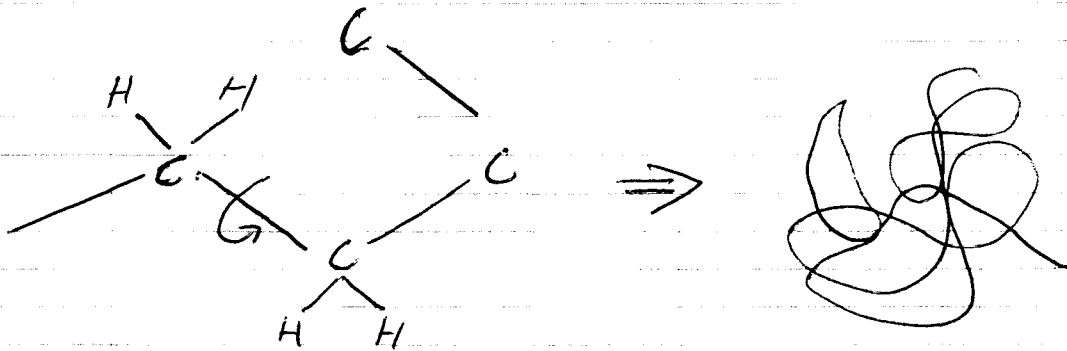
Molecular interpretation

- 1) The molecules in a system at high temperature are highly disordered. A small additional transfer of energy will result in a relatively small additional disorder. In contrast, the molecules in a system at low temperature have access to fewer energy states and the same quantity of heat will have a pronounced effect.

$$\boxed{dq_{rev} = T dS}$$

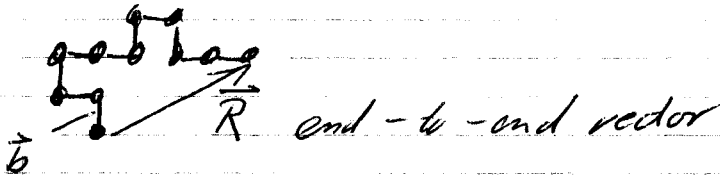
2) Entropy of a flexible polymer chain

(DeG, Introduction to Polymer Physics)



⇒ polymer chain performs a random walk since bending costs no energy!

$N$  monomers =  $N$  random steps



probability distribution of polymer end being at  $\vec{R}$ :

$$P(\vec{R}, N) = \frac{1}{Z} \sum_{i=1}^Z P(\vec{R} - \vec{b}_i, N-1)$$

↑  
possible directions  
for last polymer segment  
in 3 dimensions;  $Z=6$

$N \gg 1, |\vec{R}| \gg |b|$ :

$$P(\vec{R} - \vec{b}_i, N-1) = P(\vec{R}, N) - \frac{\partial P}{\partial N} - \sum \frac{\partial P}{\partial R_\alpha} b_{i\alpha} + \sum \frac{1}{2} \frac{\partial^2 P}{\partial R_\alpha \partial R_\beta} b_{i\alpha} b_{i\beta}$$

$$\frac{1}{Z} \sum_{i=1}^Z b_{i\alpha} = 0$$

$$\frac{1}{Z} \sum_{i=1}^Z b_{i\alpha} b_{i\beta} = \frac{\delta_{\alpha\beta} b^2}{3}$$

$$\Rightarrow \frac{\partial P}{\partial N} = \frac{b^2}{6} \frac{\partial^2 P}{\partial R^2}$$

$R=0$  for  $N=0$

$$\Rightarrow P(\vec{R}, N) = \left( \frac{3}{2\pi N b^2} \right)^{3/2} e^{-\frac{3\vec{R}^2}{2N b^2}}$$

not normalized

Gaussian distribution

as predicted by

central limit theorem

⇒ number of possibilities to achieve an end-to-end distance  $|\vec{R}|$ :

$$W(R) \sim \underbrace{z^N}_{\substack{\text{overall number of configurations} \\ \text{the polymer chain can assume}}} R^2 \cdot P(\vec{R}, N) \quad \overset{\text{Kugeloberfläche}}{=} z^N R^2 \left( \frac{3}{2\pi Nb^2} \right)^{3/2} e^{-\frac{3R^2}{2Nb^2}}$$

⇒ entropy of the polymer chain with an end-to-end distance  $R$ :

$$S_{\text{chain}} = -k_B \ln W(R) \Rightarrow \text{thermal energy stored} = -k_B T \ln W(R)$$

⇒ total energy of the polymer:

$$A_{\text{chain}} = -k_B T \ln W(R) + \text{constant} = \frac{3k_B T}{2Nb^2} R^2 + \text{constant}$$

⇒ tension pulling on the ends is

$$f = - \frac{\partial A_{\text{chain}}}{\partial R} = \frac{3k_B T}{Nb^2} R$$

⇒ Entropic force!

Entropy of phase transitions:

at the transition temperature (reversible):  $\Delta_{trs} S = \frac{\Delta_{trs} H}{T_{trs}}$

freezing, condensing  $\Rightarrow$  exothermic  $\Rightarrow \Delta_{trs} H < 0$

$\Rightarrow \Delta_{trs} S < 0$

$\Rightarrow$  system becomes more ordered

melting, vaporizing  $\Rightarrow$  endothermic  $\Rightarrow \Delta_{trs} H > 0$

$\Rightarrow \Delta S > 0$

$\Rightarrow$  system becomes more disordered

$\Rightarrow$  T Trouton's rule:

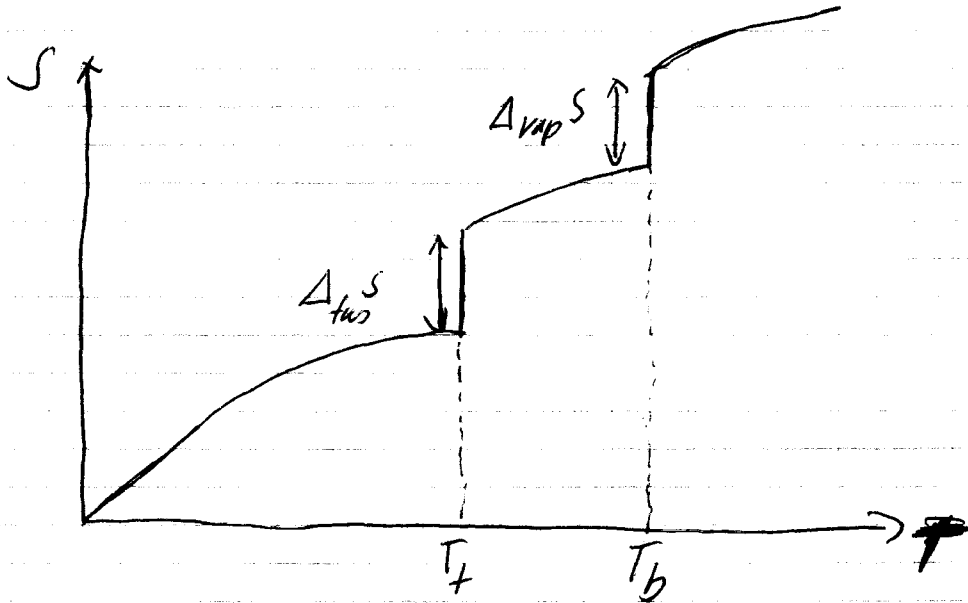
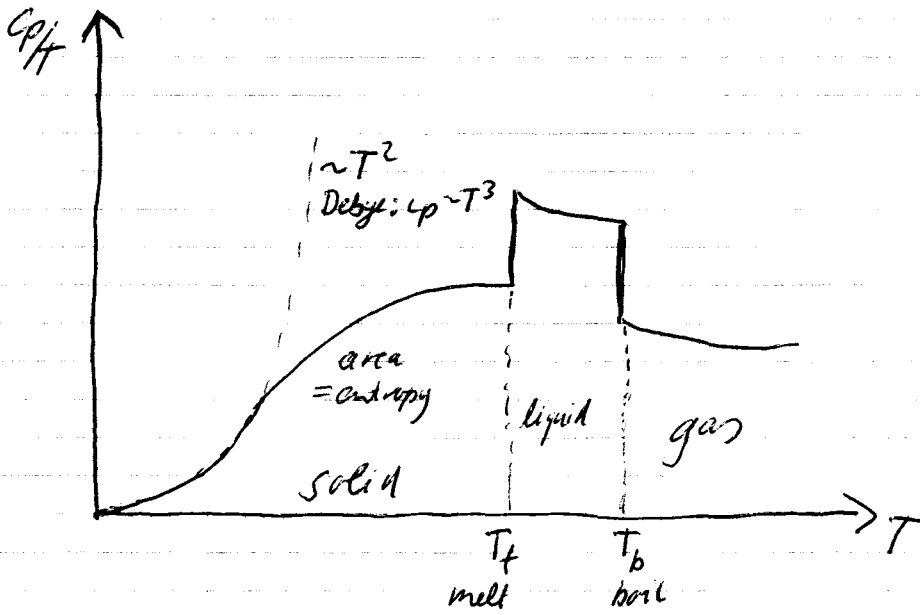
A wide range of liquids give approximately the same standard entropy of vaporization

$\approx 85 \frac{J}{K \cdot mol}$

$\Rightarrow$  comparable amount of disorder is generated when any liquid evaporates and becomes a gas

(water has a larger entropy of vaporization arising from hydrogen bonding)

Calorimetry and the measurement of entropy:



$$S(T) = S_0 + \int_0^{T_f} \frac{C_p(\text{solid})}{T} dT + \frac{\Delta_{fus} H}{T_f} + \int_{T_f}^{T_b} \frac{C_p(\text{liquid})}{T} dT + \frac{\Delta_{vap} H}{T_b} + \int_{T_b}^T \frac{C_p(\text{gas})}{T} dT$$

### 1.6. Cohesive forces:

$$V + dV: \quad U' = U + \left. \frac{\partial U}{\partial V} \right|_T dV$$

$$T + dT: \quad U' = U + \left. \frac{\partial U}{\partial T} \right|_V dT$$

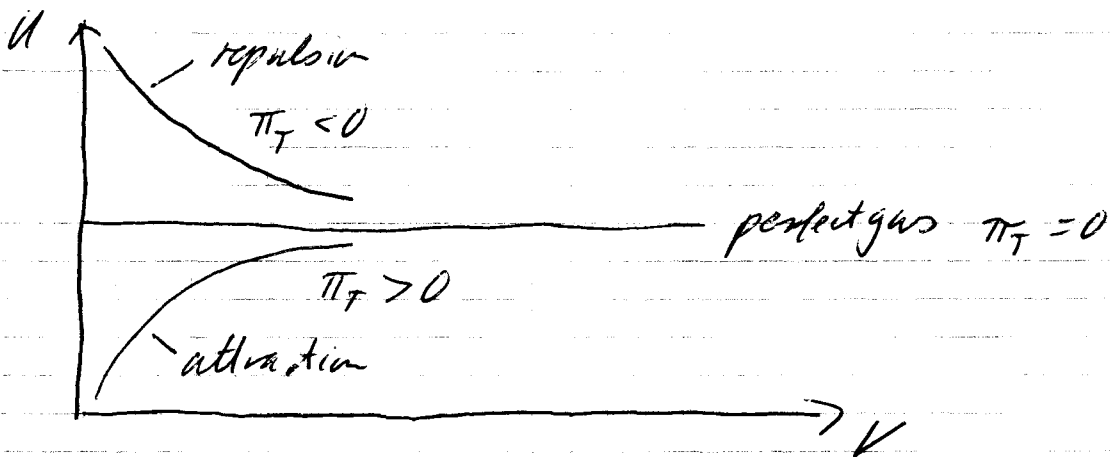
$$U' = U + \left. \frac{\partial U}{\partial V} \right|_T dV + \left. \frac{\partial U}{\partial T} \right|_V dT$$

$$dU = \left. \frac{\partial U}{\partial V} \right|_T dV + \left. \frac{\partial U}{\partial T} \right|_V dT$$

$$= \left. \frac{\partial U}{\partial V} \right|_T dV + c_V dT$$

$\left. \frac{\partial U}{\partial V} \right|_T = \pi_T$  internal pressure

$$dU = \pi_T dV + c_V dT$$



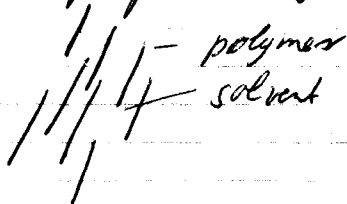
# 2. Physical Transformations of pure substances

## 2.1. Phase Diagrams

A phase of a substance is a form of matter that is uniform throughout in chemical composition and physical state (various solid phases, liquid, gas)

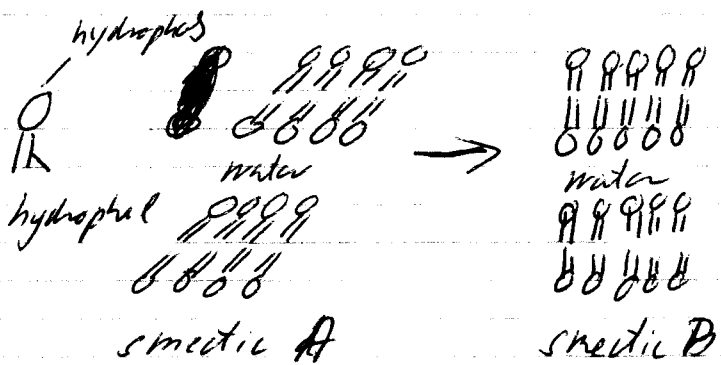
### BUT: Liquid Crystals

- rigid polymer (e.g. microtubule)



⇒ nematic phase

- Lipids ⇒ cell membranes



- chiral polymer (e.g. DNA)



cholesteric phase



⇒ Landau's definition of a phase:

phases can be distinguished by their degree of order

⇒ introduces concept of order parameter!

(Note: gas and liquid the same phase?)

phase transitions: spontaneous change of one phase into another phase

transitions can be infinitely slow

⇒ metastable states ⇒ glasses

• sand:

perfect: hexagonal order

real: jamming prevents perfect order

⇒ much less neighbors

• diamonds are metastable ⇒ ground state ~~graphite~~ graphite!

vapour pressure: pressure where liquid and vapour are at equilibrium

sublimation vapour pressure: vapour  $\leftrightarrow$  solid

boiling point, boiling temperature  $T_b$ : free vaporization throughout the liquid

critical fluid, critical temperature  $T_c$ , critical pressure  $p_c$ : liquid and gas phase are indistinguishable!

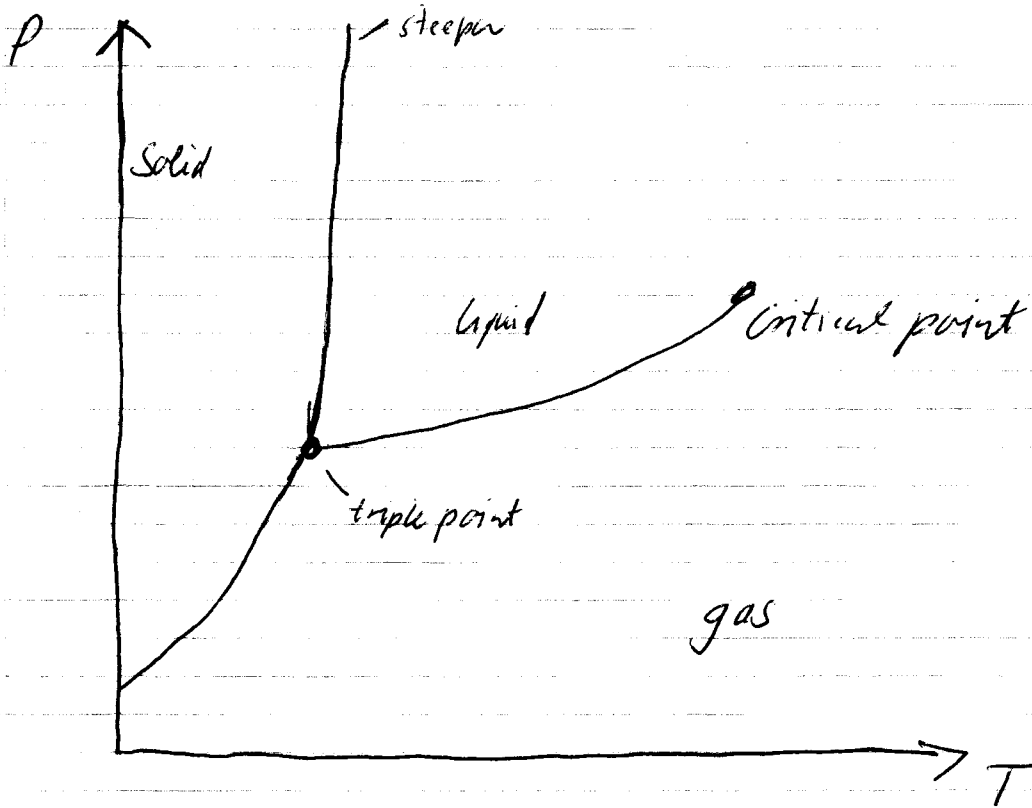
melting = freezing temperature

triple point: three phase boundaries meet

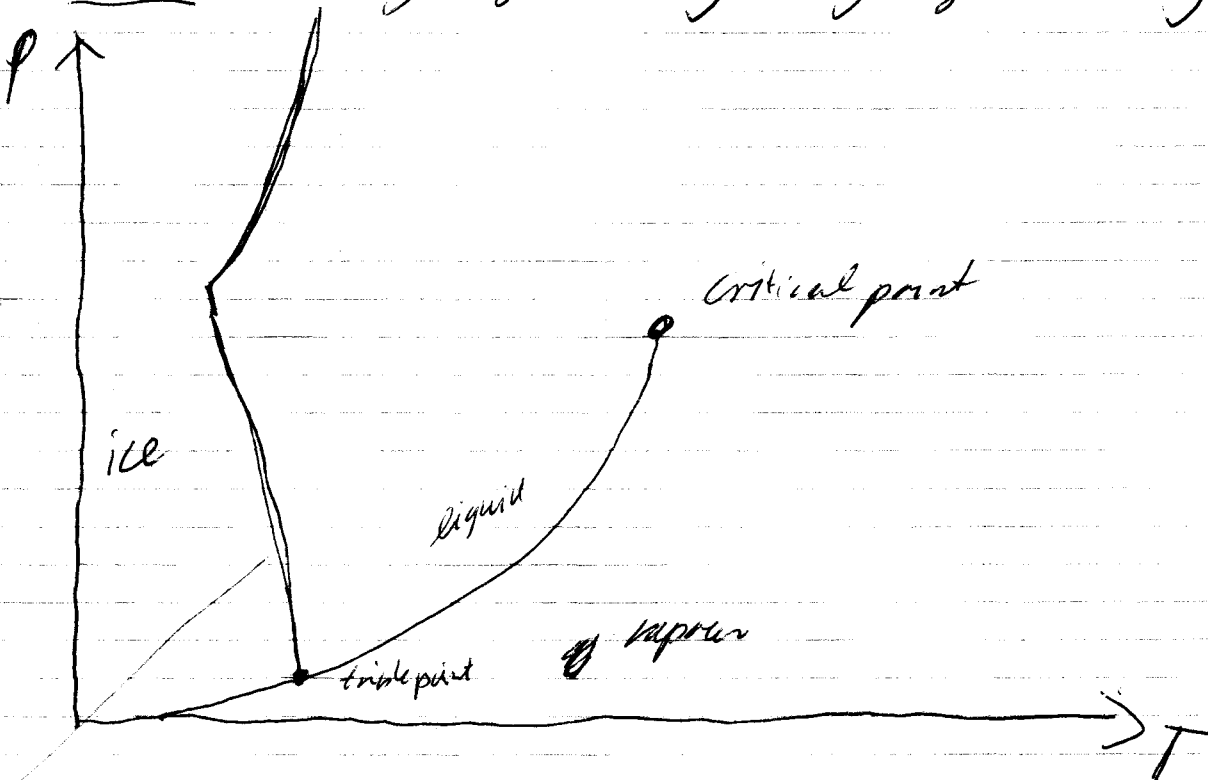
# typical phase diagrams

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1.  $\text{CO}_2$  characteristic of most substances



2. water ~~highly anomalous~~ hydrogen bonding!



Ice has larger volume than water

## 2.2. Phase stability and phase transitions

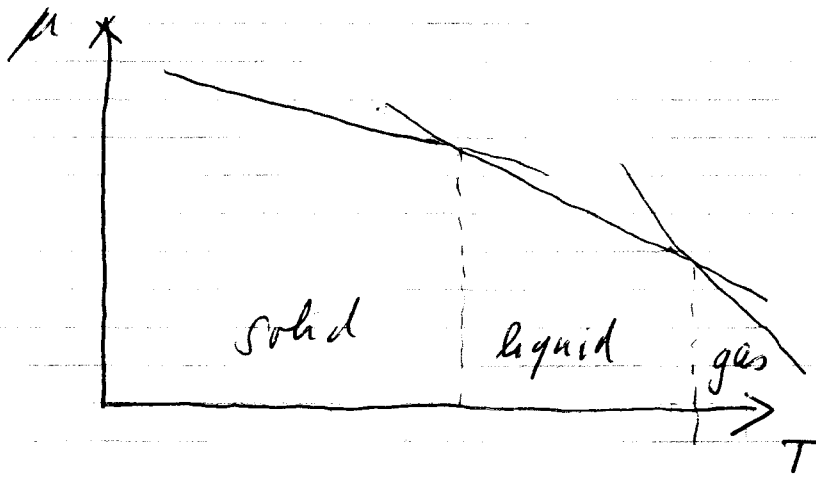
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- At equilibrium the chemical potential of a substance is the same throughout a sample regardless of how many phases are present

Proof:  $\mu_1 > \mu_2$  amount  $dn$  transferred from 1  $\rightarrow$  2  
 $\Rightarrow dG = (\mu_2 - \mu_1) dn < 0$   
spontaneous reaction

- $S(\text{gas}) > S(\text{fluid}) > S(\text{solid})$

$$\left. \frac{\partial \mu}{\partial T} \right|_P = -S$$



$$\left. \frac{\partial \mu}{\partial P} \right|_T = V \quad V(\text{liquid}) > V(\text{solid})$$

- pressure increases  $\mu$  more in a liquid than in a solid
- most substances melt at a higher temperature when subjected to pressure